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Solid State Theory I

1. Exercise Sheet

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1 ONE-PARTICLE OPERATORS IN SECOND QUANTIZATION

A one-particle operator in a many-body system is an operator that acts on the state of a single particle at a time. It is defined as $\hat{B} = \sum_{i=1}^{N} \hat{B}_i$ where \hat{B}_i is given by

$$\hat{B}_{i} = \mathbb{H}_{1} \otimes \ldots \otimes \mathbb{H}_{i-1} \otimes \sum_{\nu,\mu} \langle \nu | \hat{b} | \mu \rangle | \nu \rangle \langle \mu |_{i} \otimes \mathbb{H}_{i+1} \otimes \ldots \otimes \mathbb{H}_{N},$$
(1.1)

where $\langle \nu | \hat{b} | \mu \rangle$ are the matrix elements in one-particle quantum mechanics. In other words, \hat{B}_i acts trivially in all Hilbert spaces, except in that of particle *i*.

(a) Applying \hat{B} to the state $|\alpha_1, \alpha_2, \ldots, \alpha_N|$ introduced above, show that one obtains

$$\hat{B}|\alpha_{1},\alpha_{2},\ldots,\alpha_{N}\} = \frac{1}{\sqrt{N!}} \sum_{\mu,\nu} \langle \nu|\hat{b}|\mu\rangle$$
$$\times \sum_{i,P} \xi^{\sigma(P)} \delta_{\mu,\alpha_{P(i)}} |\alpha_{P(1)}\rangle_{1} \otimes \ldots \otimes |\nu\rangle_{i} \otimes \ldots \otimes |\alpha_{P(N)}\rangle_{N}. \quad (1.2)$$

Note that we have explicitly labeled the one-particle states with a subscript to identify the subspace they belong to within the tensor product (??), e.g. $|\nu\rangle_i$ belongs to the subspace of particle *i*.

(2 points)

(b) Show that one can write the last summation in (1.2) as

$$\sum_{i,P} \xi^{\sigma(P)} \delta_{\mu,\alpha_{P(i)}} |\alpha_{P(1)}\rangle_1 \otimes \ldots \otimes |\nu\rangle_i \otimes \ldots \otimes |\alpha_{P(N)}\rangle_N =$$
(1.3)

$$\sum_{j,P} \xi^{\sigma(P)} \delta_{\mu,\alpha_j} |\alpha_{P(1)}\rangle_1 \otimes \ldots \otimes |\nu\rangle_{P^{-1}(j)} \otimes \ldots \otimes |\alpha_{P(N)}\rangle_N.$$
(1.4)

(2 points)

Hint: Note that $\sum_{j} \delta_{j,P(i)} = 1$ for fixed *i*, as P(i) = j for some $j, \varepsilon\{1, \ldots, N\}$ since any *P* is a bijective mapping of $\{1, \ldots, N\}$ in itself.

(c) Consider again

$$\left|\alpha_{1},\ldots,\alpha_{j},\ldots,\alpha_{N}\right\rangle = \frac{1}{\sqrt{N!}}\sum_{P}\xi^{\sigma(P)}|\alpha_{P(1)}\rangle_{1}\otimes\ldots\otimes|\nu_{P(j)}\rangle_{j}\otimes\ldots\otimes|\alpha_{P(N)}\rangle_{N}.$$
(1.5)

In which position in each of these terms is α_j located? Note that I am not asking where $\alpha_{P(j)}$ is located, which is obviously at position j. From this observation and using (1.3), prove the identity

$$\frac{1}{\sqrt{N!}} \sum_{i,P} \xi^{\sigma(P)} \delta_{\mu,\alpha_{P(i)}} |\alpha_{P(1)}\rangle_1 \otimes \ldots \otimes |\nu\rangle_i \otimes \ldots \otimes |\alpha_{P(N)}\rangle_N =$$
(1.6)

$$\sum_{j} \delta_{\mu,\alpha_{j}} | \alpha_{1}, \dots, \alpha_{j-1}, \nu, \alpha_{j+1} \dots, \alpha_{N} \}.$$
(1.7)

(2 points)

(d) Show that

$$c_{\nu}^{\dagger}c_{\mu}|\alpha_{1},\ldots,\alpha_{N}\rangle = \sum_{j}\delta_{\mu,\alpha_{j}}|\alpha_{1},\ldots,\alpha_{j-1},\nu,\alpha_{j+1}\ldots,\alpha_{N}\rangle.$$
(1.8)

(2 points)

Hint: Use (??) and the symmetry properties of the wave-function under interchange of particles.

(e) Thus, show that

$$\hat{B}|\alpha_1,\ldots,\alpha_N\} = \sum_{\mu,\nu} \langle \nu|\hat{b}|\mu\rangle c_{\nu}^{\dagger}c_{\mu}|\alpha_1,\ldots,\alpha_N\}.$$
(1.9)

(1 point)

Since the state $|\alpha_1, \ldots, \alpha_N|$ is arbitrary, we conclude that a one-particle operator is given, in second quantisation, by

$$\hat{B} = \sum_{\mu,\nu} \langle \nu | \hat{b} | \mu \rangle c_{\nu}^{\dagger} c_{\mu}.$$
(1.10)

2 The density operator in second quantization

For a system of N identical particles (fermions) enclosed in a volume V with periodic boundary conditions, the density operator at position r is defined (in its first quantised form), as

$$\hat{\rho}(\boldsymbol{r}) = \sum_{i=1}^{N} \delta(\boldsymbol{r} - \hat{\boldsymbol{r}}_i), \qquad (2.1)$$

where \hat{r}_i is the position operator of particle *i*. Note that *r* is not an operator.

(a) Show that the Fourier transform $\hat{\rho}_q = \int_V d^3r e^{-iq\cdot r} \hat{\rho}(\mathbf{r})$ is given by

$$\hat{\rho}_q = \sum_{i=1}^N e^{-iq \cdot \hat{r}_i}.$$
(2.2)

(1 point)

(b) The second quantisation representation of $\hat{\rho}(\mathbf{r})$, as follows from (1.10), is given by

$$\hat{\rho}(\boldsymbol{r}) = \sum_{\sigma',\sigma} \int_{V} d^{3}x' \int_{V} d^{3}x \langle \boldsymbol{x}'\sigma' | \delta(\boldsymbol{r}-\hat{\boldsymbol{r}}) | \boldsymbol{x}\sigma \rangle \hat{\psi}_{\sigma'}^{\dagger}(\boldsymbol{x}') \hat{\psi}_{\sigma}(\boldsymbol{x}), \qquad (2.3)$$

where $\hat{\psi}_{\sigma}(\boldsymbol{x})$ and $\hat{\psi}_{\sigma}^{\dagger}(\boldsymbol{x})$ are the annihilation and creation operators for a fermion, with spin projection along z equal to σ , at point \boldsymbol{x} . Show from (2.3) that $\hat{\rho}(\boldsymbol{r})$ is given by

$$\hat{\rho}(\boldsymbol{r}) = \sum_{\sigma} \hat{\psi}^{\dagger}_{\sigma}(\boldsymbol{r}) \hat{\psi}_{\sigma}(\boldsymbol{r}).$$
(2.4)

(2 points)

(c) Show that in second quantization the FT of $\hat{\rho}(\mathbf{r})$ is given by

$$\hat{\rho}_q = \sum_{i=1}^N \hat{c}^{\dagger}_{k-q,\sigma} \hat{c}_{k,\sigma}.$$
(2.5)

(2 points)

Hint: Substitute the relations $\hat{\psi}_{\sigma}(\mathbf{r}) = \frac{1}{\sqrt{V}} \sum_{k} e^{ik \cdot r} \hat{c}_{\mathbf{k},\sigma}, \quad \hat{\psi}_{\sigma}^{\dagger}(\mathbf{r}) = \frac{1}{\sqrt{V}} \sum_{k} e^{-ik \cdot r} \hat{c}_{\mathbf{k},\sigma}^{\dagger},$ see (??), in the definition of the Fourier transform given above. Note that $\int_{V} d^{3}r e^{-i(k-k')\cdot r} = V \delta_{k,k'}.$

(d) Show that $\hat{\rho}_q^{\dagger} = \hat{\rho}_{-q}$. (1 point)

- (e) Show that $[\hat{\rho}_q, \hat{\rho}_{q'}] = 0.$ (2 points) Hint: Use the identities $[\hat{A}, \hat{B}, \hat{C}]_- = [\hat{A}, \hat{B}]_{-\xi} \hat{C} + \xi \hat{B} [\hat{A}, \hat{C}]_{-\xi}$ where $\xi = \pm 1$.
- (f) Use the first quantisation representation of ρ̂_q, as given in equation (2.2), to show the previous identity.
 (2 points)
 Hint: What is [r̂_i, r̂_j] for arbitrary i, j?
- (g) What is $\hat{\rho}_{q=0}$? (1 point)